

Meteorologische Modellierung

Introduction to Numerical Modeling
(of the Global Atmospheric Circulation)

Part II (2nd Semester)

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Outline

- 7. Elliptic Partial Differential Equations: Solving a Poisson- or Helmholtz-Equation
- 8. An example: qg barotropic channel (weather prediction)

Introduction to Numerical Modeling

7. Elliptic Partial Differential Equations: Solving a Poisson- or Helmholtz-Equation

Elliptic Partial Differential Equations: Solving a Poisson- or Helmholtz-Equation

$$a \frac{\partial^2 T}{\partial t^2} + b \frac{\partial^2 T}{\partial t \partial x} + c \frac{\partial^2 T}{\partial x^2} + d \frac{\partial T}{\partial t} + e \frac{\partial T}{\partial x} + f = 0$$

hyperbolic: $b^2 - 4ac > 0$

parabolic: $b^2 - 4ac = 0$

elliptic: $b^2 - 4ac < 0$

Elliptic equations: boundary value problem

Laplace: $\nabla^2 \Theta = 0$

Poisson: $\nabla^2 \Theta = G(x, y)$

Helmholtz: $(\nabla^2 + \lambda) \Theta = G(x, y)$

G, λ known; Θ wanted (e.g. G =vorticity; Θ =streamfunction)

Solving a Poisson-Equation: Spectral Method

$$\nabla^2 \Theta = G(x, y) \quad (1)$$

Start: Simple form of the inverse Laplace operator in spectral space (Fourier):

$$\hat{\Theta}(k, l) = \nabla^{-2} \hat{G}(k, l) = -K^{-2} \hat{G}(k, l) = -\frac{1}{k^2 + l^2} \hat{G}(k, l) \quad (2)$$

with $\hat{G}(k, l), \hat{\Theta}(k, l)$ = spectral transformed of $G(x, y)$ and $\Theta(x, y)$

k = wave number in x

l = wave number in y

Note: Trigonometric functions (Fourier, plane) or Legendere polynomials (sphere) are Eigenfunctions of the Laplace operator, i.e. $\nabla^2 E = \lambda E$

=> **Solution of (1):**

- transform $G(x, y)$ to spectral space
- compute $\hat{\Theta}(k, l)$ from (2)
- transform $\hat{\Theta}(k, l)$ to grid point space

Solving a Poisson-Equation: Grid Point Method

Start: discretization of $\nabla^2 \Theta = G(x, y)$

$y \uparrow$

	$x \rightarrow$		
	6	2	5
	3	0	1
	7	4	8

grid point nomination

For point 0:

$$(\nabla^2 \Theta)_0 - G_0 = \frac{1}{\Delta x^2} (\Theta_1 + \Theta_3 - 2\Theta_0) + \frac{1}{\Delta y^2} (\Theta_2 + \Theta_4 - 2\Theta_0) - G_0$$

!
= 0

NOTE: grid points interdependent -> no direct solution!

Solving a Poisson-Equation: Grid Point Method: Jacobi Method

Discrete Poisson-equation:

$$\nabla^2 \Theta_0 - G_0 = \frac{1}{\Delta x^2} (\Theta_1 + \Theta_3 - 2\Theta_0) + \frac{1}{\Delta y^2} (\Theta_2 + \Theta_4 - 2\Theta_0) - G_0 = 0 \quad (1)$$

Solution by Iteration (Jacobi Method):

1. choose initial field (e.g. $\Theta = 0$. or $\Theta = \text{old (known) values}$)
2. compute error ε_0 for each grid point: $\varepsilon_0 = \nabla^2 \Theta_0 - G_0$ (computed from (1))

(example: with $\Theta = 0$: $\varepsilon = G$)

3. correct each Θ_0 with error ε_0 : $\Theta'_0 = \Theta_0 + \varepsilon_0 / (\frac{2}{\Delta x^2} + \frac{2}{\Delta y^2})$

$$= (\frac{1}{\Delta x^2} (\Theta_1 + \Theta_3) + \frac{1}{\Delta y^2} (\Theta_2 + \Theta_4) - G_0) / (\frac{2}{\Delta x^2} + \frac{2}{\Delta y^2})$$

4. continue with 2. until the error is sufficiently small (given an adequate error norm)

Problem: very slow and, therefore, not feasible!

Solving a Poisson-Equation: Gauß-Seidel Method and SOR

Jacobi Method: $\Theta'_0 = (\frac{1}{\Delta x^2} (\Theta_1 + \Theta_3) + \frac{1}{\Delta y^2} (\Theta_2 + \Theta_4) - G_0) / (\frac{2}{\Delta x^2} + \frac{2}{\Delta y^2})$

Improvement: Gauß-Seidel Method

$$\Theta'_0 = (\frac{1}{\Delta x^2} (\Theta_1 + \Theta'_3) + \frac{1}{\Delta y^2} (\Theta_2 + \Theta'_4) - G_0) / (\frac{2}{\Delta x^2} + \frac{2}{\Delta y^2})$$

similar to Jacobi but using already known (computed) new values (Θ'_4, Θ'_3) advantage: feasible (fast enough)

6	2	5
3	0	1
7	4	8

grid points

More improvement: 'Successive Over-Relaxation' (SOR)

$$\begin{aligned} \Theta'_0 &= \Theta_0 + \omega \varepsilon_0 / (\frac{2}{\Delta x^2} + \frac{2}{\Delta y^2}) \\ &= \Theta_0 + \omega (\frac{1}{\Delta x^2} (\Theta_1 + \Theta'_3 - 2\Theta_0) + \frac{1}{\Delta y^2} (\Theta_2 + \Theta'_4 - 2\Theta_0) - G_0) / (\frac{2}{\Delta x^2} + \frac{2}{\Delta y^2}) \end{aligned}$$

similar to Gauß-Seidel but 'over correct' with $1 < \omega < 2$

advantage: fast disadvantage: ω needs to be chosen by 'try and error'

Solving a Poisson-Equation: Grid Point Method

Performance (NxN Grid):

Jacobi: $N^2 p/2$ iterations to decrease the initial error by factor 10^p

Gauß-Seidel: $N^2 p/4$ iterations to decrease the initial error by factor 10^p (0.5 Jacobi).

SOR: $N p/3$ iterations to decrease the initial error by factor 10^p (about $1/N$ Gauß-Seidel)

Further improvement: multigrid methods

Ansatz: Faster convergence of iteration for larger scales

=> **Multigrid method** (example):

1. interpolate $G(x,y)$ and $\Theta(x,y)$ to a coarse grid
2. solve (iterate) equation on coarse grid (e.g. by SOR)
3. interpolate solution from 2 to finer grid
4. solve (iterate) equation on finer grid (e.g. by SOR)
5. repeat 1 to 4 until final resolution (grid) and accuracy is reached

Elliptic Partial Differential Equations

Summary

- Elliptic equations: Laplace, Poisson, Helmholtz
- Spectral method: Eigenfunctions
- Grid point methods: Jacobi, Gauß-Seidel, SOR
- Multigrid method

Introduction to Numerical Modeling

8. An Example: QG Barotropic Channel Model (Weather Prediction)

The Barotropic Model

- First (functioning) numerical weather prediction model
(Charney, J. G., Fjortoft, R. and von Neumann, J. 1950.
Numerical integration of the barotropic vorticity equation.
Tellus, 2, 237–54.)
- Simple model for idealized (conceptual) studies

The Barotropic Model: Equations

Homogeneous incompressible fluid, constant density, hydrostatic equilibrium, Cartesian coordinates, z-system, barotropic, ...

Equation of motion:

$$\begin{aligned}\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} - fv &= -\frac{1}{\rho} \frac{\partial P}{\partial x} \\ \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + fu &= -\frac{1}{\rho} \frac{\partial P}{\partial y}\end{aligned}$$

Hydrostatic equation:

$$\frac{\partial P}{\partial z} = -g\rho$$

Continuity equation:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0$$

The Barotropic Model: Transformations/Approximations

1 ,primitive‘ shallow water equations

A) Integrating the hydrostatic equation and replacing the pressure gradient in the equation of motion:

$$\begin{aligned}\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} - fv &= -\frac{\partial gh}{\partial x} \\ \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + fu &= -\frac{\partial gh}{\partial y}\end{aligned}$$

B) Integrating the continuity equation (0 to h) and using boundary conditions for w (here: no bottom topography):

$$\frac{\partial gh}{\partial t} + u \frac{\partial gh}{\partial x} + v \frac{\partial gh}{\partial y} = -gh \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right)$$

The Barotropic Model: Transformations/Approximations

2 ,quasi-geostrophic‘ shallow water equations

A) QG-approximation of the equation of motion (β -plane):

$$\begin{aligned}\frac{\partial u_g}{\partial t} + u_g \frac{\partial u_g}{\partial x} + v_g \frac{\partial u_g}{\partial y} - f_0 v_a - \beta y v_g &= -\frac{\partial g h_a}{\partial x} \\ \frac{\partial v_g}{\partial t} + u_g \frac{\partial v_g}{\partial x} + v_g \frac{\partial v_g}{\partial y} + f_0 u_a + \beta y u_g &= -\frac{\partial g h_a}{\partial y}\end{aligned}$$

Or: QG vorticity equation:

$$\frac{\partial \zeta_g}{\partial t} + u_g \frac{\partial \zeta_g}{\partial x} + v_g \frac{\partial \zeta_g}{\partial y} + \beta v_g = -f_0 \left(\frac{\partial u_a}{\partial x} + \frac{\partial v_a}{\partial y} \right)$$

B) QG-approximation of the continuity equation

$$\frac{\partial g h_g}{\partial t} + u_g \frac{\partial g h_g}{\partial x} + v_g \frac{\partial g h_g}{\partial y} = -g h_0 \left(\frac{\partial u_a}{\partial x} + \frac{\partial v_a}{\partial y} \right)$$

The Barotropic Model: Transformations/Approximations

2 ,quasi-geostrophic‘ shallow water equations

with geostrophic streamfunction $\psi = g h_g / f_0$:

$$\begin{aligned}u_g &= -\frac{\partial \Psi}{\partial y} \\ v_g &= \frac{\partial \Psi}{\partial x} \\ \zeta_g &= \nabla^2 \Psi = \frac{\partial^2 \Psi}{\partial x^2} + \frac{\partial^2 \Psi}{\partial y^2}\end{aligned}$$

$$\begin{aligned}\frac{\partial (\nabla^2 - \lambda^{-2}) \Psi}{\partial t} &= -\frac{\partial \Psi}{\partial x} \frac{\partial (\nabla^2 - \lambda^{-2}) \Psi}{\partial y} + \frac{\partial \Psi}{\partial y} \frac{\partial (\nabla^2 - \lambda^{-2}) \Psi}{\partial x} - \beta \frac{\partial \Psi}{\partial x} \\ &= -J(\Psi, (\nabla^2 - \lambda^{-2}) \Psi) - \beta \frac{\partial \Psi}{\partial x}\end{aligned}$$

with λ = Rossby Radius of Deformation = $(g h_0)^{1/2} / f_0$

The Barotropic Model: Transformations/Approximations

3 non-divergent shallow water equations

(QG-)vorticity equation (non divergent):

$$\frac{\partial \zeta_g}{\partial t} + u_g \frac{\partial \zeta_g}{\partial x} + v_g \frac{\partial \zeta_g}{\partial y} + \beta v_g = 0$$

Or:

$$\begin{aligned} \frac{\partial \nabla^2 \Psi}{\partial t} &= -\frac{\partial \Psi}{\partial x} \frac{\partial \nabla^2 \Psi}{\partial y} + \frac{\partial \Psi}{\partial y} \frac{\partial \nabla^2 \Psi}{\partial x} - \beta \frac{\partial \Psi}{\partial x} \\ &= -J(\Psi, \nabla^2 \Psi) - \beta \frac{\partial \Psi}{\partial x} \end{aligned}$$

The Barotropic Model: Transformations/Approximations

4 equivalent barotropic model

Assumption: absolute value of velocity may change with height but not the direction: $(u, v) = A(p)(<u>(x, y, t), <v>(x, y, t))$

$$ = \frac{1}{p_s} \int_0^{p_s} B dp$$

It follows:

$$\frac{\partial (\nabla^2 - A(p_s) \lambda^{-2}) \Psi^*}{\partial t} = -J(\Psi^*, \nabla^2 \Psi^*) - \beta \frac{\partial \Psi^*}{\partial x}$$

with $\Psi^* = <A^2> \Psi$; $<A> = 1$

Valid for the equivalent barotropic level p^* with: $A(p^*) = <A^2>$

(typically 600-500hPa; minimum divergence)

Summary

The Barotropic Model: Equations

,primitive' equations:

$$\begin{aligned}\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} - fv &= -\frac{\partial gh}{\partial x} \\ \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + fu &= -\frac{\partial gh}{\partial y} \\ \frac{\partial gh}{\partial t} + u \frac{\partial gh}{\partial x} + v \frac{\partial gh}{\partial y} &= -gh \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right)\end{aligned}$$

Quasi-Geostrophic: $\frac{\partial(\nabla^2 - \lambda^{-2})\Psi}{\partial t} = -J(\Psi, (\nabla^2 - \lambda^{-2})\Psi) - \beta \frac{\partial \Psi}{\partial x}$

Non-divergent: $\frac{\partial \nabla^2 \Psi}{\partial t} = -J(\Psi, \nabla^2 \Psi) - \beta \frac{\partial \Psi}{\partial x}$

Equivalent barotropic: $\frac{\partial(\nabla^2 - A(p_s)\lambda^{-2})\Psi^*}{\partial t} = -J(\Psi^*, \nabla^2 \Psi^*) - \beta \frac{\partial \Psi^*}{\partial x}$

From Equations to Numerical Model:

Model Design

A) The equation(s): $\frac{\partial \nabla^2 \Psi}{\partial t} = -J(\Psi, \nabla^2 \Psi) - \beta \frac{\partial \Psi}{\partial x}$
Here: barotropic non-divergent

B) The numerical method - general
Here: grid point method

C) The numerical method – specific
variables, operators, grid, discretizations, work flow, boundary conditions, etc.

Barotropic non-divergent Model: Numerics

$$\frac{\partial \nabla^2 \Psi}{\partial t} = -J(\Psi, \nabla^2 \Psi) - \beta \frac{\partial \Psi}{\partial x} \Leftrightarrow \frac{\partial \Psi}{\partial t} = \nabla^{-2} \left[-J(\Psi, \nabla^2 \Psi) - \beta \frac{\partial \Psi}{\partial x} \right]$$

Prognostic variable: Streamfunction Ψ

Operators:

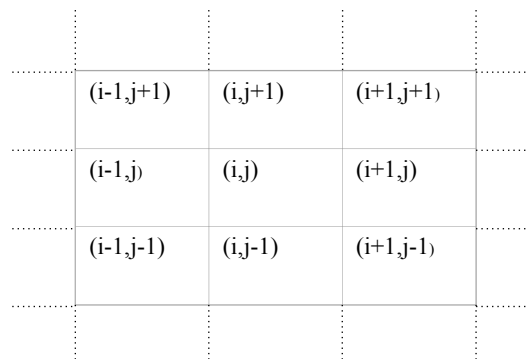
- Derivation in time: $\frac{\partial \Psi}{\partial t}$
- Derivation in space: $\beta \frac{\partial \Psi}{\partial x}$
- Jacobi-Operator: $J(\Psi, \nabla^2 \Psi)$
- Laplace-Operator and its inverse: $\Psi \leftrightarrow \nabla^2 \Psi$

Barotropic non-divergent Model: Numerics

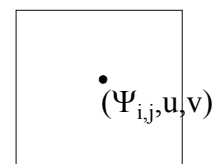
The Grid:

Lon-Lat grid; one prognostic variable only (Ψ) -> Arakawa A

The Grid:



Gridbox (i,j)



Barotropic non-divergent Model: Numerics

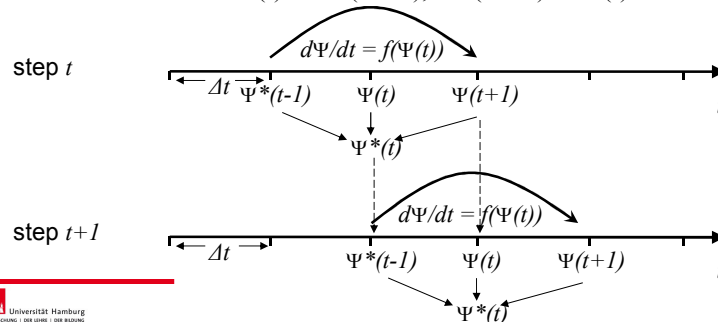
Discretizations

Time derivative: Leapfrog with Robert-Asselin filter

(Three level scheme; conservative (non-dissipative) system;
Computation of $\Psi(t)$ (and $\Psi(t+\Delta t)$) by weighted averages of
 $\Psi(t+\Delta t)$, $\Psi(t)$ and $\Psi(t-\Delta t)$)

Courant-Friedrich-Levy criterion: $(u\Delta t/\Delta x)^2 < (1-\gamma)^2$

- Calculation rule:
1. $\Psi(t+\Delta t) = \Psi^*(t-\Delta t) + 2\Delta t \cdot f(\Psi(t))$
 2. $\Psi^*(t) = \Psi(t) + \gamma(\Psi^*(t-\Delta t) - 2\Psi(t) + \Psi(t+\Delta t))$ ($\gamma = \text{filter const.}$)
 3. $\Psi^*(t) \rightarrow \Psi^*(t-\Delta t)$; $\Psi(t+\Delta t) \rightarrow \Psi(t)$ (e.g. $\gamma = 0.1$)



Barotropic non-divergent Model: Numerics

Discretizations

Derivation in space: centered differences

$$\frac{\partial \Psi}{\partial x} = \frac{1}{2\Delta x} (\Psi_1 - \Psi_3)$$

$$\frac{\partial \Psi}{\partial y} = \frac{1}{2\Delta y} (\Psi_2 - \Psi_4)$$

Grid:

6 = (i-1, j+1)	2 = (i, j+1)	5 = (i+1, j+1)
3 = (i-1, j)	0 = (i, j)	1 = (i+1, j)
7 = (i-1, j-1)	4 = (i, j-1)	8 = (i+1, j-1)

Laplacian: centered differences

$$\nabla^2 \Psi(i, j) = \nabla^2 \Psi_0 = \frac{1}{\Delta x^2} (\Psi_1 + \Psi_3 - 2\Psi_0) + \frac{1}{\Delta y^2} (\Psi_2 + \Psi_4 - 2\Psi_0)$$

Barotropic non-divergent Model: Numerics Discretizations

Jacobi-Operator: $J(\Psi, \nabla^2 \Psi)$

$$\begin{aligned} \text{analytically } J(a, b) &= \frac{\partial a}{\partial x} \frac{\partial b}{\partial y} - \frac{\partial a}{\partial y} \frac{\partial b}{\partial x} \\ &= \frac{\partial}{\partial x} \left(a \frac{\partial b}{\partial y} \right) - \frac{\partial}{\partial y} \left(a \frac{\partial b}{\partial x} \right) \\ &= \frac{\partial}{\partial y} \left(b \frac{\partial a}{\partial x} \right) - \frac{\partial}{\partial x} \left(b \frac{\partial a}{\partial y} \right) \end{aligned}$$

Grid:

6 = (i-1, j+1)	2 = (i, j+1)	5 = (i+1, j+1)
3 = (i-1, j)	0 = (i, j)	1 = (i+1, j)
7 = (i-1, j-1)	4 = (i, j-1)	8 = (i+1, j-1)

$$\begin{aligned} \text{numerically: } J_1 &= \frac{1}{4\Delta x \Delta y} \{ (a_1 - a_3)(b_2 - b_4) - (a_2 - a_4)(b_1 - b_3) \} \\ J_2 &= \frac{1}{4\Delta x \Delta y} \{ a_1(b_5 - b_8) - a_3(b_6 - b_7) - a_2(b_5 - b_6) + a_4(b_8 - b_7) \} \\ J_3 &= \frac{1}{4\Delta x \Delta y} \{ b_2(a_5 - a_6) - b_4(a_8 - a_7) - b_1(a_5 - a_8) + b_3(a_6 - a_7) \} \end{aligned}$$

$J = (J_1 + J_2 + J_3)/3$ (Arakawa '66) enstrophy and energy conserving

Barotropic non-divergent Model: Numerics Discretizations

Inverse Laplacian (solution of a Poisson-equation): $\nabla^2 \Theta = G(x, y)$

e.g. 'Successive Over-Relaxation' (SOR)

Discrete Laplacian:

$$\nabla^2 \Theta_0 - G_0 = \frac{1}{\Delta x^2} (\Theta_1 + \Theta_3 - 2\Theta_0) + \frac{1}{\Delta y^2} (\Theta_2 + \Theta_4 - 2\Theta_0) - G_0 = 0$$

Grid:

6 = (i-1, j+1)	2 = (i, j+1)	5 = (i+1, j+1)
3 = (i-1, j)	0 = (i, j)	1 = (i+1, j)
7 = (i-1, j-1)	4 = (i, j-1)	8 = (i+1, j-1)

Iterative solution:

1. estimate the error: $\varepsilon_0 = \nabla^2 \Theta_0 - G_0$

2. (over-) correct the error (using already new values):

$$\Theta'_0 = \Theta_0 + \omega \varepsilon_0 / \left(\frac{2}{\Delta x^2} + \frac{2}{\Delta y^2} \right)$$

$$= \Theta_0 + \omega \left(\frac{1}{\Delta x^2} (\Theta_1 + \Theta'_3 - 2\Theta_0) + \frac{1}{\Delta y^2} (\Theta_2 + \Theta'_4 - 2\Theta_0) - G_0 \right) / \left(\frac{2}{\Delta x^2} + \frac{2}{\Delta y^2} \right)$$

Barotropic non-divergent Model: Numerics

Work flow:
$$\frac{\partial \nabla^2 \Psi}{\partial t} = -J(\Psi, \nabla^2 \Psi) - \beta \frac{\partial \Psi}{\partial x}$$

1. Initialization

- a) define grid
- b) set boundary conditions
- c) set initial conditions

2. Time loop

- a) compute the right hand side (tendency)
- b) Inverse Laplacian = solving a Poisson-equation
- c) compute Ψ at new time step

3. Finalization

- a) write restart files

Summary: Numerics

Equation: *barotropic non-divergent*
$$\frac{\partial \nabla^2 \Psi}{\partial t} = -J(\Psi, \nabla^2 \Psi) - \beta \frac{\partial \Psi}{\partial x}$$

Method: Grid point

Grid: Arakawa A

Prognostic variable: streamfunction

Time stepping: Leapfrog with filter

Differentials in space: central differences

Jacobi operator: energy and enstrophy conserving (Arakawa)

Inverse Laplacian: 'Successive Over-Relaxation' (SOR)

From the Design to the Code: The FORTRAN program

Modular structure:

(sets of) subroutines for individual model parts:

- input
- output
- grid definition
- derivation in space
- Laplacian
- Jacobi operator
- Leapfrog
- Euler
- 'Successive Over-Relaxation' (SOR)
- boundary conditions
- initial conditions
- ...

organized and called by a main program

From the Design to the Code: The FORTRAN program

How to program?

- use 'modules' for global parameters/variables
- use a name convention (real/integer; global/local, etc)
- document/comment your code (the more the better)
- try to be flexible (use parameters etc.)
- test as frequent as possible

From the Design to the Code: The FORTRAN program

How to test?

- use simple structures with know solutions (sin,cos) to check the derivatives
- use analytic solution to check the dynamics e.g. Rossby-Haurwitz wave

R-H wave:

only one wave number k and l in x- and y-direction: $\Psi(x, y, t) = \Psi_0 e^{i(kx+ly-\omega t)}$

$$\Rightarrow J(\Psi, \nabla^2 \Psi) = 0$$

$$\Rightarrow \frac{\partial \nabla^2 \Psi}{\partial t} = -\beta \frac{\partial \Psi}{\partial x}$$

$$\Rightarrow c = -\frac{\beta}{K^2}$$

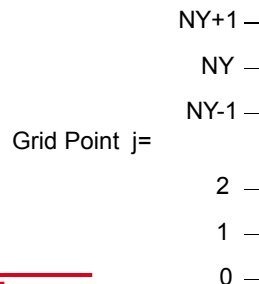
**westward propagating, amplitude and
wavenumber conserving**

The Barotropic Model: Boundary Conditions

General: at each boundary (East, West, North, South) two boundary conditions are needed: u and v or streamfunction Ψ and vorticity ξ

Our model: non divergent barotropic vorticity equation, grid point method, A-grid, Ψ

The grid: Example: one longitude:



Here: $0, NY+1$ = Channel boundaries, $1, 2, \dots, NY$ = 'aktive'
i.e. boundary conditions at GP 0 and $NY+1$ needed

Boundary Conditions: Examples

$NY+1$
 NY
 1
 0

a) Prescribed from data

b) Cyclic: $A_0=A_{NY}$ and $A_{NY+1}=A_1$ (with $A=u, v$ or Ψ, ξ)

c) No flow across boundary: $v_0=v_{NY+1}=0$, i.e. $\Psi_0=\Psi_{NY+1}=\text{const}$ in x , as $v_{0,NY+1}=\left(\frac{\partial \Psi}{\partial x}\right)_{0,NY+1}$

and

c1) $u_0=u_1$ and $u_{NY+1}=u_{NY}$ (**full slip boundary condition**), i.e. $\xi_{0,NY+1}=0$,

as $\xi_{0,NY+1}=-\left(\frac{\partial u}{\partial y}\right)_{0,NY+1} \stackrel{!}{=} 0$

c2) $u_{0,NY+1}=0$ (**no slip boundary condition**), i.e. (in our model) $\xi_{0,NY+1}=2\frac{\Psi_{1,NY+1}-\Psi_{0,NY}}{(\Delta y)^2}$

as $\xi_j=-\frac{u_{j+1/2}-u_{j-1/2}}{\Delta y}$ and $u_{-1/2,NY+3/2} \stackrel{!}{=} u_{1/2,NY+1/2}$ and $u_{1/2,NY+1/2}=-\frac{\Psi_{1,NY+1}-\Psi_{0,NY}}{\Delta y}$

Boundary Conditions: Summary

- Prescribed
- Cyclic
- Perpendicular: No flow across boundary
- Tangential: No slip and full slip